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OPTICAL PROPERTIES OF HgGaInS₄ CRYSTALS

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Abstract. Luminescence, transmission, reflection, and wavelength modulation transmission and wavelength modulation reflection spectra of HgGalnS₄ crystals were researched in the temperature range of 300-10 K. This article discusses the various features observed in the spectra of HgGalnS₄ crystals. The dependence on temperature of the energy gap is shown. Spectra of wavelength modulation transmission have minima a₁, a₂, and a₃, which vary with temperature and are caused by indirect transitions L- Γ with phonon absorption and emission. Spectra of thicker crystals have more features caused by the same indirect transitions. The photoluminescence show a strong maximum at 2.43 eV, attributed to radiative recombination of free excitons. These various features have significant implications for the understanding of the electronic and optical properties of HgGalnS₄ crystals.

Keywords: *Photouminescence, transmission, reflection, wavelength modulation transmission and reflection, absorption, emission.*

Rezumat. Au fost cercetate spectrele de luminescență, de transmisie, de reflexie și modulate în funcție de lungimea de undă a transmisiei și reflexiei ale cristalelor HgGaInS₄ în intervalul de temperatură 300-10 K. În acest articol se discută diferite caracteristici observate în spectrele cristalelor HgGaInS₄. Este prezentată dependența de temperatură a energiei de bandă. Spectrele de modificare a transmisiei în funcție de lungimea undei, cu minime a₁, a₂ și a₃, care variază cu temperatura și sunt cauzate de tranziții indirecte L-Γ cu absorbție și emisie de fononi. Spectrele cristalelor mai groase au mai multe caracteristici cauzate de aceleași tranziții indirecte. Spectrele de fotoluminescență arată un maxim puternic la 2.43 eV, atribuit recombinării radiative a excitonilor liberi. Aceste diferite caracteristici au implicări semnificative pentru înțelegerea proprietăților electronice și optice ale cristalelor HgGaInS₄.

Cuvinte cheie: Luminescență, transmisie, reflexie, transmisie și reflexie modulate în funcție de lungimea de undă, absorbție, emisie.

1. Introduction

Crystals HgGalnS₄ formed when defect chalcopyrite HgGa₂S₄ interacts with spinel HgIn₂S₄ and has layered structure. The lattice structure of HgGalnS₄ crystals consists of an arrangement of anions in a hexagonal close-packed structure, with layers of cations occupying particular arrangement both voids in the shape of tetrahedrons and octahedrons [1].

The hexagonal lattice of HgGaInS₄ crystals has unit cell parameters (lengths of the cell edges) of a = b = 0.39 nm and c = 3.14 nm [1]. HgGaInS₄ crystals exhibit similar chemical, crystalline and electrophysical properties compared to $ZnIn_2S_4$, CdGaInS₄, CdIn₂S₂Se₂ crystals [2].

The highly pronounced layered structure of $ZnIn_2S_4$ crystals, which belong to the $A^{"}B_2^{""}C_4^{\vee i}$ group of ternary semiconducting compounds, makes them particularly interesting and has attracted the attention of researchers, leading to numerous experimental studies [3-4]. These crystals possess photo-sensitivity and exhibit photoluminescence [5-6], showcasing intriguing optical properties.

The classic thiospinel structure is the crystalline form of $HgIn_2S_4$ compounds. The optical properties of compounds in this group have been predominantly studied using crystals like $ZnIn_2S_4$, $ZnAl_2Se_4$, and $CdGa_2S_4$ [7-13]. These crystals are utilized in various optical filters and nonlinear optical devices. Investigated materials such as $HgIn_2S_4$ and $HgGa_2Se_4$ are intriguing due to their linear and nonlinear optical characteristics, as well as their narrow band gap of 1.2-1.6 eV [9,14].

2. Materials and Methods

The single crystals of HgGaInS₄ were grown by the gas transport method in ampoules. The material was grown with mirror-like surfaces with dimensions of ~10x7 mm². The plates had different thicknesses (0.01-3 mm) and were easily cleavable. Monocrystals of nanoscale thickness (50 nm) were visually distinguishable by their color. The spectra at low temperatures of the material placed in a helium cryostat LTS-22 C 330 cryogenic optical system were obtained using an MDR-2 spectrometer with an aperture with a ratio of 1:2 and a linear dispersion of 0.7 nm/mm. The optical system of the MDR-2 spectrometer is shown in Figure 1.



Figure 1. Schematic representation of set-up based on single spectrometer MDR-2 for transmission and reflection spectra registration.

The photoluminescence spectra were measured using the double-beam spectrophotometer SDL-1 with a light amplification ratio of 1:2 and a linear dispersion of 0.7 nm/mm. The optical system of SDL-1 is shown in Figure 2.



Figure 2. Scheme of measurement installation based on double spectrometer SDL-1 for photoluminescence spectra registration.

3. Results

Initial results on the absorption and photoluminescence of HgGaInS₄ crystals have been previously investigated [2]. The energy gap of HgGaInS₄ is equal to 2.41 eV at temperature 300 K and 2.46 eV at 80 K [2]. It was also noted that the exponential tail of the

absorption edge is caused by the existence of states that are distributed quasi-continuously [15].

The newly discovered layered phase of HgGaInS₄ has a crystalline structure that includes empty spaces or vacancies within the cation lattice and the atoms in the cation sublattice are randomly distributed. This material is an electron semiconductor, with a band gap of 2.41 eV, high sensitivity light, strong to photoluminescence. In the band exponentially qap, there are distributed electron traps (with a characteristic energy of about 65 meV) near the conduction band bottom, and several distinct energy levels positioned roughly at 0.01, 0.11, and 0.2 eV above the



Figure 3. Absorption edge of HgGalnS₄ crystals, with thicknesses 0.1-0.6 μ m within the temperature interval of 300-10 K and with a thickness of 1.2 μ m (insert).

top of the valence band. The initial two levels are associated with radiative electron transitions from donor-like states that are exponentially distributed. The third level is a sensitizing level. The slope of electron transitions from the valence band to quasi-

continuously distributed states in the exponential region of the absorption edge remains constant regardless of temperature.

Figure 3 shows the spectra of absorption of HgGaInS₄ thicknesses of crystals 0.1-0.6 μ m at different temperatures in the interval of 300-10 K. Decreasing temperature leads to the absorption edge shifts towards higher energies. Measurements were conducted also on crystals with thicknesses of 1-1.2 μ m in the range where absorption coefficients are higher than 10³ cm⁻¹ (see insert from Figure 3). Maxima can be detected on the spectra at energies of 2.42-2.53 eV.



Figure 4. Temperature dependence of wavelength modulation transmission ($\Delta T/\Delta \lambda$) (a) and of a₁, a₂ and a₃ minima positions (b).

The spectra of the wavelength modulation transmission change $\Delta T/\Delta \lambda$ of HgGalnS₄ crystals at temperatures of 10-300K exhibit features (minima) a₁, a₂, and a₃, the energy positions of which vary with temperature, as shown in Figure 4a and Figure 4b, positions of peaks is shown in table 1.

| Temperature values and pe | Temperature values and peak positions used in the research | |
|---------------------------|--|---|
| Temperature, K | Peak position, eV | _ |
| 10 | a ₁ – 2.188 | |
| | a ₂ – 2.231 | |
| | a ₃ - 2.261 | |
| 20 | a ₁ – 2.191 | |
| | a ₂ – 2.242 | |
| | a ₃ - 2.269 | |
| 50 | a ₁ – 2.232 | |
| | a ₂ – 2.273 | |
| | a ₃ – 2.303 | |
| 100 | a ₁ – 2.222 | |
| | a ₂ – 2.267 | |
| | a ₃ – 2.295 | |
| 200 | a ₁ – 2.193 | |
| | a ₂ – 2.241 | |

| Continuation Table | |
|------------------------|-----|
| a3 – 2.269 | |
| a1 – 2.175 | 250 |
| a ₂ – 2.224 | |
| a3 – 2.252 | |
| a1 – 2.161 | 300 |
| a2 – 2.206 | |
| a3 – 2.237 | |
| | |



Figure 5. Spectra of wavelength modulation transmission ($\Delta T/\Delta \lambda$) for HgGalnS₄ crystals with a thickness of 600µm at temperatures of 20 and 50K.



Figure 6. Spectra of absorption (K) of HgGalnS₄ with different thicknesses (right side – 0.5mm, left side - 0.1mm), at temperatures of 10 K and 300 K.

On the other hand, massive crystals with thickness about (600 μ m) exhibit minima x1-x5, as shown in Figure 5. We assume that x1-x5 and a₁-a₃ are caused by indirect transitions L- Γ (E_g^{ind}-2.2-2.3 eV) with phonon absorption and emission. Thicker crystals and additional investigations are required for a more accurate determination of this interval.

Figure 6 shows the shift of the absorption edge spectra of HgGalnS₄ crystals with a thickness of 0.5 mm in the temperature interval of 300-10 K. As seen from the spectra, at temperature 300K the absorption edge starts at 1.9 eV, at temperature 10 K, at 2.08 eV. At an absorption coefficient of 10^3 cm⁻¹, the temperature coefficient of the absorption edge shift (Δ Eg/ Δ T) equal to 8.1·10⁻⁴ eV/K. Decreasing temperature leads to a shift of the absorption edge towards higher energies. At absorption coefficients of 3·10³ cm⁻¹, coefficient of the temperature shift of the absorption edge (Δ Eg/ Δ T) is 6·10⁻⁴ eV/K, and at an absorption coefficient of 3·10⁴ cm⁻¹, θ (Δ Eg/ Δ T) is twice as small and equal to -3·10⁻⁴ eV/K. These data practically confirm the concept that changes in temperature have a weaker effect on deep energy levels and high-energy electronic transitions.

Figure 7 shows the reflection spectra of 0.5 mm thick HgGalnS₄ crystals at temperature 10 K in the high absorption region. In the energy interval of 2.4-2.6 eV, a weak change in reflection spectra is observed with a maximum at 2.4 eV. The change in intensity from the maximum to the minimum occurs within 6-7%. This characteristic is more clearly distinguished in the reflection spectra wavelength modulation. The observed R feature is due to direct electronic transitions into the exciton zone at the Γ point.

Measured excitonic reflection ($R_{exper.}$) and calculated using a one-oscillator model ($R_{calc.}$) spectra is presented in Figure 5.



Figure 7. Reflection spectra and spectra of wavelength modulation reflection ($\Delta R/\Delta \lambda$) of HgGaInS₄ crystals with a thickness of 0.5mm at temperature 10K.

The resonant value of the excitonic transition is 2.43 eV, the background dielectric constant is 9.83, the splitting between longitudinal and transverse components is 13 meV, and the factor of damping is about 187. Naturally, with such a large damping factor, the excitonic reflection spectra contour is weakly distinguished.



Figure 8. Absorption spectra of HgGaInS₄ crystals with a thickness of 0.1 μ m and photoluminescence spectra with excitation wavelength 325 nm laser at 10 K.

Absorption spectra of crystals with a thickness of 0.1 μ m clearly show peaks at E₁ (2.450 eV), E₂ (2.580 eV), and E₂[•] (2.605 eV), as shown in Figure 8. The photoluminescence spectral dependence of HgGalnS₄ crystals with excitation wavelength 325 nm at temperatures of 20 K and 10 K shows an intense peak at 2.439 eV, as can be seen Figure 8. The comparison of the absorption and photoluminescence peaks suggests that the E₁ absorption peak and E_{PL} photoluminescence peak are associated with the excitonic transition. The E₂ and E₂[•] absorption peaks are associated with the n=2 and n=3 states, respectively. In this case, the exciton binding energy (Rydberg constant) is 180 meV, and the bandgap width is 2.619 eV.

The interference fringes are clearly visible in the spectra, and as can be seen, the interference spectra extend more into the short-wavelength region as the crystals become thinner. The refractive index was determined by using the relation between the energy positions of the observed maxima (or minima) and the refractive index:

$$\boldsymbol{n} = \frac{M}{2d\left(\frac{1}{\lambda_1} - \frac{1}{\lambda_2}\right)} \quad , \tag{1}$$

M equal to 1 for nearby maxima (minima); λ_1 and λ_2 position value (nm) for maxima (minima) in interference spectra.



Figure 9. Transmission spectra of crystals with different thicknesses (a-270 μm, b-120 μm, c-0.1 μm) at 10 K (a). Dependence of the refractive index n(interf.) calculated from interference transmission spectra (T) and from the Kramers-Kronig calculation of the reflection spectra (n-Kram-Kr) (b).

The spectral dependence of the refractive index $n_{interf.}$ obtained from the calculation of interference transmission spectra (T) is presented in Figure 9b. The obtained spectra are compared with the refractive index spectral dependence $n_{Kram-Kr.}$ obtained through calculations of the reflection spectra R applying the Kramers-Kronig equation. The spectral dependences of the reflection coefficient (R) and the optical constants refractive index (n), extinction coefficient (k), real part of permittivity (ϵ_1), imaginary part of permittivity (ϵ_2), and phase (ϕ) are interconnected by Kramers-Kronig relationships. These relationships allow the calculation of one parameter from another and are useful in analyzing the optical properties of materials.

$$\begin{cases} r = \frac{N-1}{N+1} = \frac{n+ik-1}{n+ik+1} \\ r = \sqrt{R}e^{-i\varphi} = \sqrt{R}\left(\cos\varphi - i\sin\varphi\right) \end{cases} \begin{cases} n = \frac{1-R}{1-2\sqrt{R}\cos\varphi + R} \\ k = \frac{2\sqrt{R}\sin\varphi}{1-2\sqrt{R}\cos\varphi + R} \end{cases}$$
(2)

As can be seen, the refractive index values and the spectral dependence obtained by the two methods are in good agreement.

4. Discussion

The study of the HgGaInS₄ crystals revealed interesting properties related to their optical absorption, reflection, and transmission spectra. The absorption edge of the material shifts towards higher energies as the temperature decreases, indicating that the temperature has a lesser effect on deep energy levels and high-energy electronic transitions [1]. The thickness of the crystal affects the appearance of features in the absorption spectra, which may be caused by indirect transitions with phonon absorption and emission [16].

The reflection spectra of the crystals exhibit a weak change in the intensity in the energy interval of 2.4-2.6 eV, with a maximum at 2.43 eV, which is attributed to direct electronic transitions into the exciton zone at the Γ point. The photoluminescence spectra also show peaks related to exciton modes [16].

The transmission spectra exhibit interference fringes that extend more into the shortwavelength region as the crystals become thinner, and the refractive index values obtained by the two methods (refractive index obtained from the calculation of interference spectra and obtained through calculations applying Kramers-Kronig equation) are in good agreement.

Overall, the study of the HgGaInS₄ crystals provided valuable insights into the optical properties of the material, which could be useful for applications in optoelectronics and photonics. Further investigations on the material's properties, particularly on thicker crystals, are required for a more accurate determination of the observed features and transitions.

5. Conclusions

The study of HgGalnS₄ crystals' spectra has revealed several significant features. The shift in the absorption edge towards higher energies with decreasing temperature indicates a temperature dependence of the bandgap energy. The wavelength modulation transmission spectra exhibit energy minima caused by indirect transitions L- Γ with phonon absorption and emission, and the spectra of thicker crystals have additional minima due to the same transitions. The weak change in the reflection spectra at energies between 2.4-2.6 eV, with a maximum at 2.4 eV, is resulted by direct electronic transitions to excitonic zone at the Γ point. Spectra of photoluminescence with an excitation by a 325 nm show a strong maximum at 2.21 eV, resulting from the radiative recombination of free excitons. The absorption edge temperature shift coefficient varies with the absorption coefficient, indicating that temperature changes affect deep energy levels and high-energy electronic transitions to a lesser extent. Overall, these results provide valuable insights into the electronic and optical properties of HgGalnS₄ crystals.

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Conflicts of Interest: The author declares no conflict of interest.

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